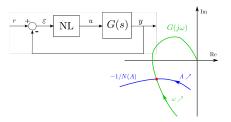
4AESE - Analyse des Systèmes Non-Linéaires

## Chapitre 4 : Describing functions

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version : 11-2022



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## Introduction

**2** Harmonic linearization

Self-oscillations



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#### Introduction

Describing functions method is an extension of harmonic method for some nonlinearities in a closed-loop system

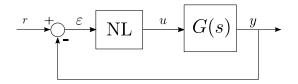
↔ in french, it is called : méthode du premier harmonique

- It approximates a nonlinear element by a "equivalent" linear term
  - ↔ harmonic linearization
- Method particularly used to predict limit cycle in a closed-loop system



#### Framework

In this chapter, we only consider closed-loop system of the form



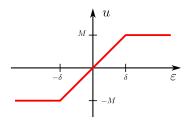
Assumptions

- The nonlinear element NL is a separable term
- It is time-invariant
- The linear term, G(s), is stable and a low-pass filter type (known as filtering hypothesis)



#### Nonlinear element : example 1

#### Saturation

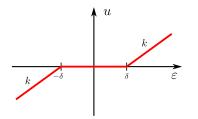


- Linear for  $\varepsilon \in [-\delta, \delta]$ , constant for large values of  $|\varepsilon|$ .
- Often models actuator limitations
  - $\hookrightarrow$  power amplifiers, motors, servo-value for flow control
- Usually caused by limits on component size, properties of materials, available power, mechanical configuration...



#### Nonlinear element : example 2

Dead zone

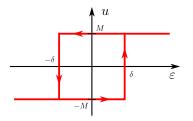


- *u* is zero until the magnitude of the input exceeds some threshold  $|\varepsilon| > \delta$ .
- Usually characterize actuators (valve, motor...) that are unresponsive to low input signals
- For instance, it models static friction on motor shaft



#### Nonlinear element : example 3

#### Hysteresis



# Examples 1 and 2 are static nonlinearities, also named memoryless

 $\hookrightarrow$  the output solely depends on the instantaneous input value

Hysteresis output depends on the instantaneous and past input values

→ nonlinear element with *memory* 



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#### **Frequency response**

**Recall for linear case** : The response to a sine function is also a sine function (at steady state)

 $\hookrightarrow$  with the same frequency  $\omega$  but different amplitude and phase shift w.r.t.  $\omega$ 

$$\begin{array}{c|c} u(s) & y(s) \\ u(t) = u_0 \sin(\omega t) & y(t) = y_0 \sin(\omega t + \varphi) \end{array}$$

For nonlinear case : The response is a periodic signal (at steady state)

 $\hookrightarrow y(t) = y(t + T)$ , then a Fourier series expansion can be used

$$u(t) \longrightarrow \mathbf{NL} \qquad y(t) \longrightarrow \mathbf{y}(t) = y(t+T)$$



#### First harmonic approximation

In the closed-loop system, let's assume there is a limit cycle and the oscillating signal is

$$arepsilon(t)=A\sin(\omega t)$$
 (  $\equiv -y(t)$  if reference is 0 )

Fourier series expansion of the nonlinear component response

$$u(t) = a_0 + \sum_{n=1}^{\infty} a_n \cos(n\omega t) + b_n \sin(n\omega t)$$

- Fourier coefficients  $a_0$ ,  $a_n$  and  $b_n$  are functions of A and  $\omega$
- if the nonlinearity is odd,  $a_0 = 0$  (often the case)



#### First harmonic approximation

The *filtering hypothesis* implies all harmonics are filtered out and only the fundamental component is considered :

 $u(t) \simeq a_1 \cos(\omega t) + b_1 \sin(\omega t)$   $\simeq M \sin(\omega t + \phi)$  $\simeq M e^{j(\omega t + \phi)}$ 

with

$$a_{1}(A,\omega) = \frac{\omega}{\pi} \int_{(T)} u(t) \cos(\omega t) dt \qquad b_{1}(A,\omega) = \frac{\omega}{\pi} \int_{(T)} u(t) \sin(\omega t) dt$$
$$M(A,\omega) = \sqrt{a_{1}^{2} + b_{1}^{2}} \qquad \phi(A,\omega) = \arctan(\frac{a_{1}}{b_{1}})$$



#### **Describing function**

Similarly to linear case, frequency response = ratio sinusoidal output / sinusoidal input

The *describing function* of a nonlinear element is the complex ratio of the fundamental component of the nonlinearity by the input sinusoid

$$\mathsf{N}(A,\omega) = rac{Me^{j(\omega t+\phi)}}{Ae^{j\omega t}} = rac{M}{A}e^{j\phi} = rac{1}{A}(b_1+ja_1)$$



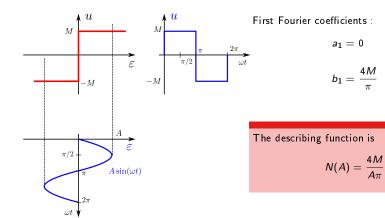
The approximated frequency response depends on the input amplitude A

← this operation is called quasi-linearization

 $\blacktriangleright$  For static nonlinearities, the describing function is independent of  $\omega$ 

#### Example 1 : relay





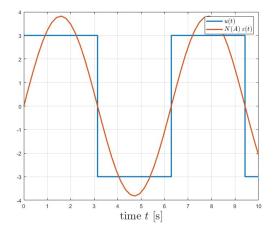
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#### Example 1 : relay

Simulation for :  $\omega = 1$ , A = 2 and M = 3

 $\hookrightarrow$  equivalent gain N(A) = 1.909

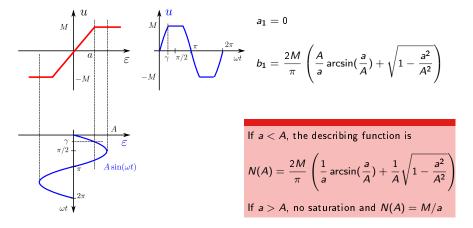


 $\Diamond$  Keeping in mind that u is then filtered by a low-pass type transfer function

#### Example 2 : saturation



First Fourier coefficients :

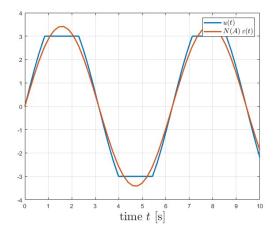




#### Example 2 : saturation

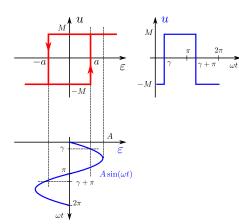
Simulation for :  $\omega = 1$ , A = 2, M = 3 and a = 1.5

 $\hookrightarrow$  equivalent gain N(A) = 1.711



 $\Diamond$  Keeping in mind that u is then filtered by a low-pass type transfer function

#### **Example 3 : hysteresis**



First Fourier coefficients

$$a_{1} = -\frac{4M}{\pi}\frac{a}{A}$$
$$b_{1} = \frac{4M}{\pi}\sqrt{1-\frac{a^{2}}{A^{2}}}$$

If a < A, the describing function is  $N(A) = \frac{4M}{A\pi} \left( \sqrt{1 - \frac{a^2}{A^2}} - j \frac{a}{A} \right)$ 

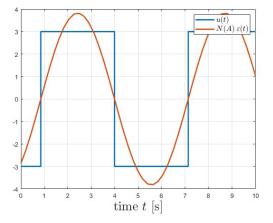




#### **Example 3 : hysteresis**

Simulation for :  $\omega = 1$ , A = 2, M = 3 and a = 1.5

$$\hookrightarrow$$
 equivalent gain  $N(A) = 1.26 - 1.43i$ , soit  $\left\{ egin{array}{c} M = 1.909 \\ \phi = -0.848 \end{array} 
ight.$ 



 $\Diamond$  Keeping in mind that u is then filtered by a low-pass type transfer function



#### **Computing describing functions**

Different methods to compute a describing function of a nonlinear element

 $u = f(\varepsilon)$ 

- Analytical calculation. when the nonlinear characteristic is known, explicit and simple enough to calculate the integrals (a<sub>1</sub> and b<sub>1</sub>); the nonlinearity could also be approximated by piecewise linear functions. Result is an analytical expression of N(A, ω).
- Numerical integration. when the nonlinear characteristic is given by a graph / table of values, integrals numerically computed with a discrete sums of surface over small intervals (numerical algorithm). Result is a plot of N w.r.t A and ω.
- Experimental evaluation. Interesting when no information about the nonlinearity (or too complex), but can be isolated and excited with sinusoidal inputs for various A and ω; compute the ratio of amplitude and phase shift with the output. Result is a plot of N w.r.t A and ω.

Chapitre 4 : Describing functions └─Self-oscillations



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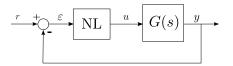
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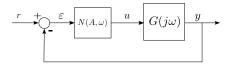


#### **Closed-loop system analysis**

Let's go back to closed-loop system (with r = 0)



Assuming the system is oscillating, the closed-loop can be approximated by



 $\Diamond$  The output must satisfy the relationship :  $y = G(j\omega)N(A,\omega)(-y)$ 



#### **Existence of limit cycles**

It implies that

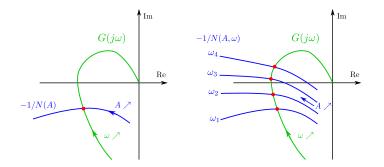
$$G(j\omega)N(A,\omega)+1=0 \qquad \Leftrightarrow \qquad G(j\omega)=-rac{1}{N(A,\omega)}$$

- If some solutions exist, there is (are) limit cycle(s) with amplitude A and frequency \u03c6 (approximately)
- if not, there is no limit cycle
- lt is 2 nonlinear equations with 2 variables (A and  $\omega$ )
- May be very difficult to solve analytically for high-order systems
- Usually, graphical approach
  - $\hookrightarrow$  plot  $G(j\omega)$  and  $-1/N(A,\omega)$  in the complex plane
  - $\hookrightarrow$  find the intersection points



#### Existence of limit cycles : graphical method

Illustration of the method when the describing function depends only on the amplitude A (left) and on both amplitude A and frequency  $\omega$  (right)



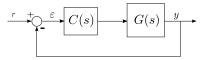
Note that in both cases there could be several intersection points



#### Stability of limit cycles

Previous slides were about **detecting existence** of limit cycles. What about their **stability**?

Before that, let's recall the Nyquist criterion



Characteristic equ. of the closed-loop system : 1 + C(s)G(s) = 0 (or CG = -1)

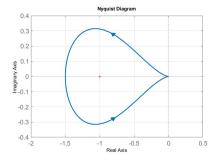
#### Nyquist criterion

Procedure :

- **b** Draw in the complex plane C(s)G(s), s following the Nyquist path
- Determine the number N of clockwise encirclement around the point (-1, 0)
- Determine the number P of unstable poles of C(s)G(s)
- Then Z = N + P is the number of unstable poles of the closed-loop system



Example : with 
$$C(s) = 1$$
 and  $G(s) = \frac{3}{(s+2)(s-1)}$ 



• Number of clockwise encirclement around the point (-1,0) : N=-1

Number of unstable poles of C(s)G(s) : P = 1

• Then, there is Z = N + P = 0 unstable pole for the closed-loop system

 $\hookrightarrow$  closed-loop system stable

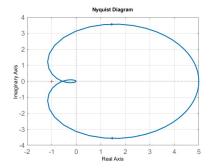


Simplified version : when C(s)G(s) has no unstable pole

(critère de Revers)

 $\Rightarrow$  no encirclement around the point (-1,0)  $\rightarrow$  closed-loop system stable

Example with 
$$C(s)G(s)=rac{10}{(s+1)^2(s+2)}$$



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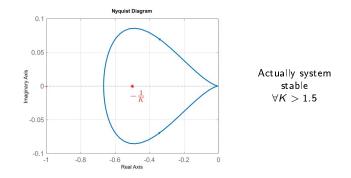


**Simple extension** : as a function of a tunable gain K in the loop

The characteristic equation : 1 + KC(s)G(s) = 0 or  $C(s)G(s) = -\frac{1}{K}$ 

 $\Rightarrow$  check the encirclement around the point (-1/K, 0)

Example with 
$$C(s)G(s)=rac{1}{(s+1.5)(s-1)}$$
 and  $K=2$ 

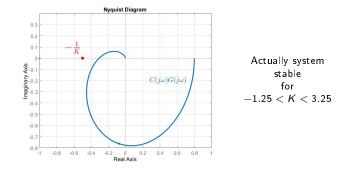


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Another Example with 
$$C(s)G(s) = \frac{1}{s^3 + 2s^2 + 2.25s + 1.25}$$
 and  $K = 2$ 

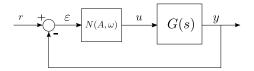
 $(poles = -1 and -0.5 \pm j)$ 





#### Stability of limit cycles

Back to closed-loop system with a describing function :



Characteristic equation  $1 + N(A, \omega)G(s) = 0$  or  $G(s) = -\frac{1}{N(A, \omega)}$ 

• Check encirclement around the point 
$$\left(R_e[-\frac{1}{N}], I_m[-\frac{1}{N}]\right)$$

▶ By assumption G(s) is stable  $\rightarrow$  no unstable pole

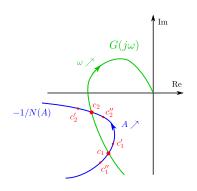
• Check if critical points is on left or right of  $G(j\omega)$  locus when  $\omega \nearrow$ 

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#### **Graphical analysis**

Consider a system where two limit cycles are predicted



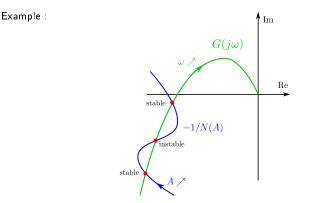
- At c<sub>1</sub>, there is a limit cycle with an amplitude A<sub>1</sub> and a frequency ω<sub>1</sub>
  - if a slight disturbance increases A; we move to c'<sub>1</sub>; the system is unstable; the amplitude continues to increase; we move along curve -1/N(A) toward c<sub>2</sub>
  - if a slight disturbance decreases A; we move to c<sub>1</sub>"; the system is stable; the amplitude continues to decrease; we move along curve -1/N(A) toward 0
  - the limit cycle is unstable
- At c<sub>2</sub>, there is a limit cycle with an amplitude A<sub>2</sub> and a frequency ω<sub>2</sub>
  - if a slight disturbance increases A; we move to  $c_2^\prime$ ; the system is stable; the amplitude decreases; we move back toward  $c_2$
  - if a slight disturbance decreases A; we move to c<sub>2</sub><sup>''</sup>; the system is unstable; the amplitude increases; we move back toward c<sub>2</sub>
  - the limit cycle is stable



#### Stability condition (graphical)

#### Loeb criterion

A limit cycle of amplitude  $A_0$  and frequency  $\omega_0$  is stable if the intersection point is such that along the nyquist plot of  $G(j\omega)$  as  $\omega$  increases, the direction of increasing A along the critical curve -1/N(A) is toward the left.

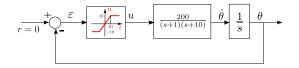


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#### Example

Simple control of a DC motor with a saturation



What is the describing function of the nonlinearity?

- Show that a limit cycle exists.
- What would be the approximated amplitude and frequency of the self-oscillations?
- Is the limit cycle stable?

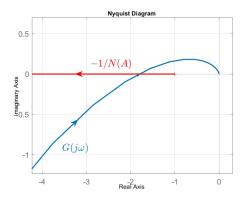
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#### Example

• Describing function : 
$$N(A) = \frac{2}{\pi} \arcsin\left(\frac{10}{A}\right) + \frac{20}{A\pi}\sqrt{1 - \frac{100}{A^2}}$$
 (if  $A > 10$ )

▶ Nyquist plot of  $G(j\omega)$  and -1/N(A); note that N(10) = 1 and  $N(+\infty) = 0$ 



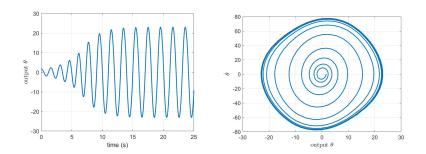
From the plot,  $\omega = 3.22 \text{ rad/s} (T = 1.95 \text{ s})$  and A = 22.1

The limit cycle is stable

#### Example



Simulation of the closed loop : ouput (left) and phase plane (right)

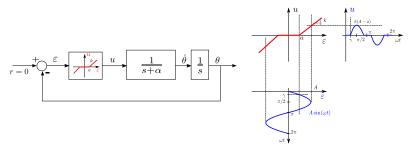


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#### Exercise

Simple control of a DC motor with a dead zone



What is the describing function of the nonlinearity?

- > Draw a sketch of the nyquist plot of  $G(j\omega)$  and the critical locus -1/N(A).
- Does a limit cycle exists?



#### Solution

