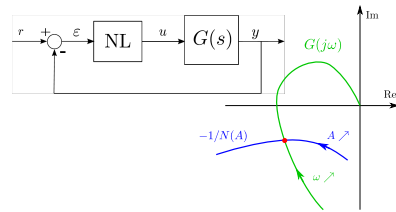


Chapitre 4 : Describing functions

Yassine ARIBA



Sommaire

- 1 Introduction
- 2 Harmonic linearization
- 3 Self-oscillations

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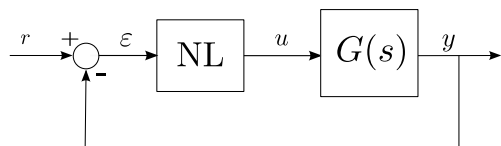
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Introduction

- ▶ Describing functions method is an extension of harmonic method for some nonlinearities in a closed-loop system
 - ↔ in french, it is called : *méthode du premier harmonique*
- ▶ It approximates a nonlinear element by a "equivalent" linear term
 - ↔ harmonic linearization
- ▶ Method particularly used to predict limit cycle in a closed-loop system

Framework

In this chapter, we only consider closed-loop system of the form

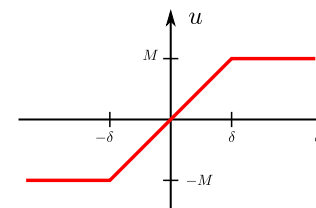


Assumptions :

- ▶ The nonlinear element NL is a separable term
- ▶ It is time-invariant
- ▶ The linear term, $G(s)$, is stable and a low-pass filter type (known as *filtering hypothesis*)

Nonlinear element : example 1

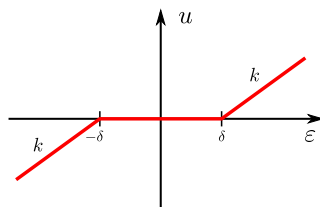
Saturation



- ▶ Linear for $\varepsilon \in [-\delta, \delta]$, constant for large values of $|\varepsilon|$.
- ▶ Often models actuator limitations
 - ↪ power amplifiers, motors, servo-valve for flow control
- ▶ Usually caused by limits on component size, properties of materials, available power, mechanical configuration...

Nonlinear element : example 2

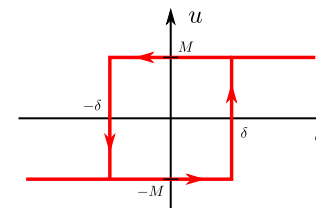
Dead zone



- ▶ u is zero until the magnitude of the input exceeds some threshold $|\varepsilon| > \delta$.
- ▶ Usually characterize actuators (valve, motor...) that are unresponsive to low input signals
- ▶ For instance, it models static friction on motor shaft

Nonlinear element : example 3

Hysteresis



- ▶ ...
- ▶ Examples 1 and 2 are static nonlinearities, also named *memoryless*
 - ↪ the output solely depends on the instantaneous input value
- ▶ Hysteresis output depends on the instantaneous and past input values
 - ↪ nonlinear element with *memory*

Sommaire

1 Introduction

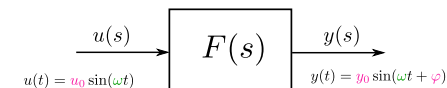
2 Harmonic linearization

3 Self-oscillations

Frequency response

Recall for linear case : The response to a sine function is also a sine function (at steady state)

↔ with the same frequency ω but different amplitude and phase shift w.r.t. ω



For nonlinear case : The response is a periodic signal (at steady state)

↔ $y(t) = y(t + T)$, then a Fourier series expansion can be used



First harmonic approximation



In the closed-loop system, let's assume there is a limit cycle and the oscillating signal is

$$\varepsilon(t) = A \sin(\omega t) \quad (\equiv -y(t) \text{ if reference is } 0)$$

Fourier series expansion of the nonlinear component response

$$u(t) = a_0 + \sum_{n=1}^{\infty} a_n \cos(n\omega t) + b_n \sin(n\omega t)$$

- ▶ Fourier coefficients a_0 , a_n and b_n are functions of A and ω
- ▶ if the nonlinearity is odd, $a_0 = 0$ (often the case)

First harmonic approximation

The *filtering hypothesis* implies all harmonics are filtered out and only the fundamental component is considered :

$$u(t) \simeq a_1 \cos(\omega t) + b_1 \sin(\omega t)$$

$$\simeq M \sin(\omega t + \phi)$$

$$\simeq M e^{j(\omega t + \phi)}$$

with

$$a_1(A, \omega) = \frac{\omega}{\pi} \int_{(T)} u(t) \cos(\omega t) dt \quad b_1(A, \omega) = \frac{\omega}{\pi} \int_{(T)} u(t) \sin(\omega t) dt$$

$$M(A, \omega) = \sqrt{a_1^2 + b_1^2} \quad \phi(A, \omega) = \arctan\left(\frac{a_1}{b_1}\right)$$

Describing function

Similarly to linear case, frequency response = ratio sinusoidal output / sinusoidal input

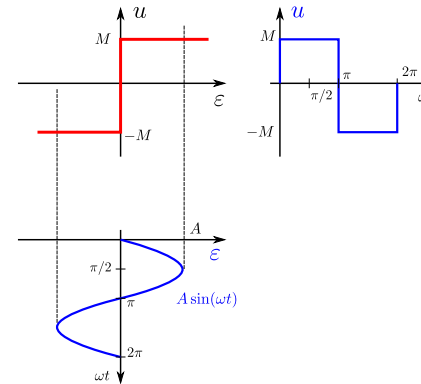
The *describing function* of a nonlinear element is the complex ratio of the fundamental component of the nonlinearity by the input sinusoid

$$N(A, \omega) = \frac{M e^{j(\omega t + \phi)}}{A e^{j\omega t}} = \frac{M}{A} e^{j\phi} = \frac{1}{A} (b_1 + j a_1)$$



- ▶ The approximated frequency response depends on the input amplitude A
 - ↳ this operation is called quasi-linearization
- ▶ For static nonlinearities, the describing function is independent of ω

Example 1 : relay



First Fourier coefficients :

$$a_1 = 0$$

$$b_1 = \frac{4M}{\pi}$$

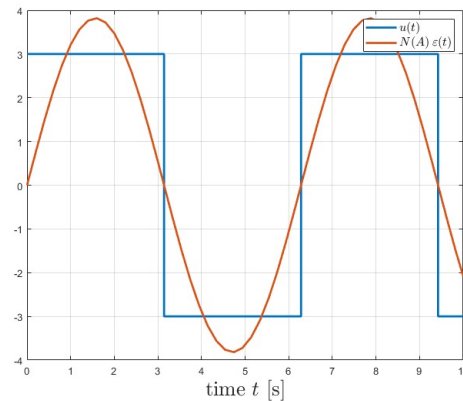
The describing function is

$$N(A) = \frac{4M}{A\pi}$$

Example 1 : relay

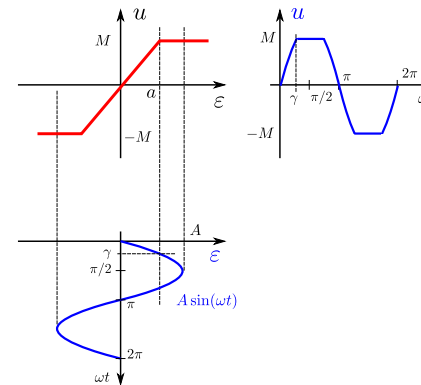
Simulation for : $\omega = 1$, $A = 2$ and $M = 3$

↳ equivalent gain $N(A) = 1.909$



◇ Keeping in mind that u is then filtered by a low-pass type transfer function

Example 2 : saturation



First Fourier coefficients :

$$a_1 = 0$$

$$b_1 = \frac{2M}{\pi} \left(\frac{A}{a} \arcsin\left(\frac{a}{A}\right) + \sqrt{1 - \frac{a^2}{A^2}} \right)$$

If $a < A$, the describing function is

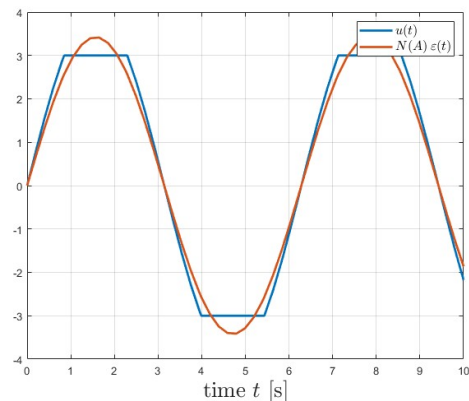
$$N(A) = \frac{2M}{\pi} \left(\frac{1}{a} \arcsin\left(\frac{a}{A}\right) + \frac{1}{A} \sqrt{1 - \frac{a^2}{A^2}} \right)$$

If $a > A$, no saturation and $N(A) = M/a$

Example 2 : saturation

Simulation for : $\omega = 1$, $A = 2$, $M = 3$ and $a = 1.5$

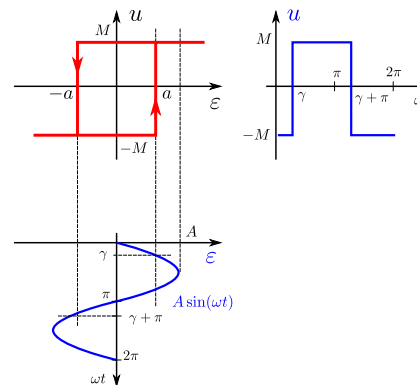
↳ equivalent gain $N(A) = 1.711$



◇ Keeping in mind that u is then filtered by a low-pass type transfer function

Example 3 : hysteresis

First Fourier coefficients :



$$a_1 = -\frac{4M}{\pi} \frac{a}{A}$$

$$b_1 = \frac{4M}{\pi} \sqrt{1 - \frac{a^2}{A^2}}$$

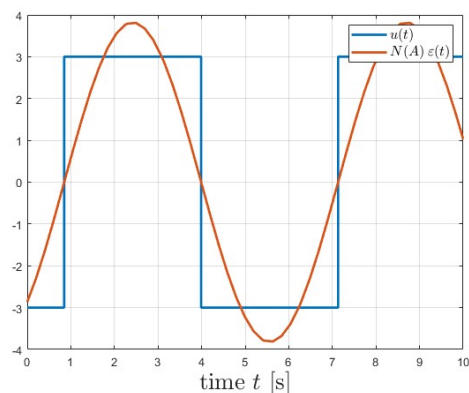
If $a < A$, the describing function is

$$N(A) = \frac{4M}{A\pi} \left(\sqrt{1 - \frac{a^2}{A^2}} - j \frac{a}{A} \right)$$

Example 3 : hysteresis

Simulation for : $\omega = 1$, $A = 2$, $M = 3$ and $a = 1.5$

↳ equivalent gain $N(A) = 1.26 - 1.43j$, soit $\begin{cases} M = 1.909 \\ \phi = -0.848 \end{cases}$



◇ Keeping in mind that u is then filtered by a low-pass type transfer function

Computing describing functions

Different methods to compute a describing function of a nonlinear element

$$u = f(\varepsilon)$$

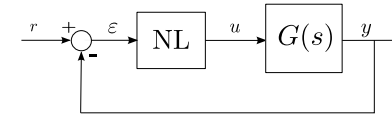
- ▶ **Analytical calculation.** when the nonlinear characteristic is known, explicit and simple enough to calculate the integrals (a_1 and b_1) ; the nonlinearity could also be approximated by piecewise linear functions. Result is an analytical expression of $N(A, \omega)$.
- ▶ **Numerical integration.** when the nonlinear characteristic is given by a graph / table of values, integrals numerically computed with a discrete sums of surface over small intervals (numerical algorithm). Result is a plot of N w.r.t A and ω .
- ▶ **Experimental evaluation.** Interesting when no information about the nonlinearity (or too complex), but can be isolated and excited with sinusoidal inputs for various A and ω ; compute the ratio of amplitude and phase shift with the output. Result is a plot of N w.r.t A and ω .

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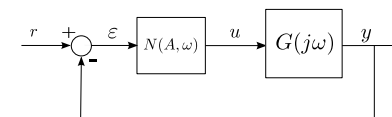
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Closed-loop system analysis

Let's go back to closed-loop system (with $r = 0$)



Assuming the system is oscillating, the closed-loop can be approximated by



◇ The output must satisfy the relationship : $y = G(j\omega)N(A, \omega)(-y)$

Existence of limit cycles

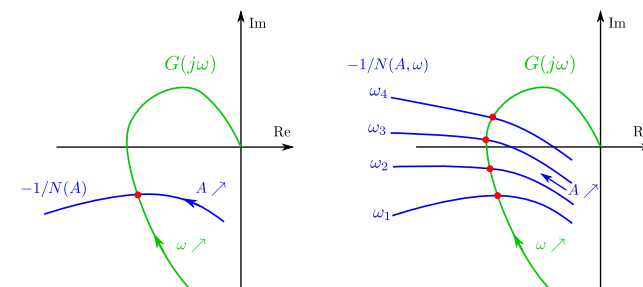
It implies that

$$G(j\omega)N(A, \omega) + 1 = 0 \quad \Leftrightarrow \quad G(j\omega) = -\frac{1}{N(A, \omega)}$$

- ▶ If some solutions exist, there is (are) limit cycle(s) with amplitude A and frequency ω (approximately)
- ▶ if not, there is no limit cycle
- ▶ It is 2 nonlinear equations with 2 variables (A and ω)
- ▶ May be very difficult to solve analytically for high-order systems
- ▶ Usually, graphical approach
 - ↪ plot $G(j\omega)$ and $-1/N(A, \omega)$ in the complex plane
 - ↪ find the intersection points

Existence of limit cycles : graphical method

Illustration of the method when the describing function depends only on the amplitude A (left) and on both amplitude A and frequency ω (right)

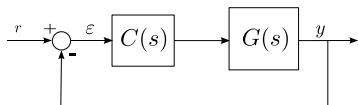


Note that in both cases there could be several intersection points

Stability of limit cycles

Previous slides were about **detecting existence** of limit cycles. What about their **stability**?

Before that, let's recall the Nyquist criterion



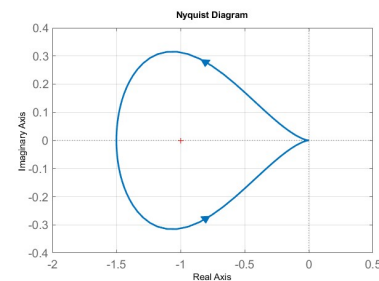
Characteristic equ. of the closed-loop system : $1 + C(s)G(s) = 0$ (or $CG = -1$)

Nyquist criterion

Procedure :

- ▶ Draw in the complex plane $C(s)G(s)$, s following the Nyquist path
- ▶ Determine the number N of clockwise encirclement around the point $(-1, 0)$
- ▶ Determine the number P of unstable poles of $C(s)G(s)$
- ▶ Then $Z = N + P$ is the number of unstable poles of the closed-loop system

Example : with $C(s) = 1$ and $G(s) = \frac{3}{(s+2)(s-1)}$



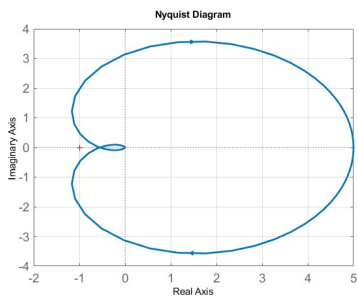
- ▶ Number of clockwise encirclement around the point $(-1, 0)$: $N = -1$
- ▶ Number of unstable poles of $C(s)G(s)$: $P = 1$
- ▶ Then, there is $Z = N + P = 0$ unstable pole for the closed-loop system
↔ closed-loop system stable

Simplified version : when $C(s)G(s)$ has no unstable pole

(critère de Revers)

⇒ no encirclement around the point $(-1, 0)$ → closed-loop system stable

Example with $C(s)G(s) = \frac{10}{(s+1)^2(s+2)}$

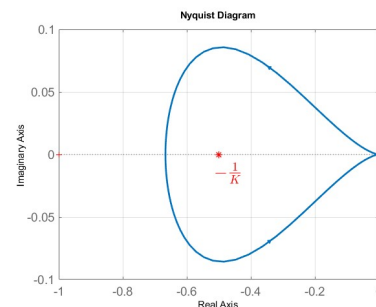


Simple extension : as a function of a tunable gain K in the loop

The characteristic equation : $1 + KC(s)G(s) = 0$ or $C(s)G(s) = -\frac{1}{K}$

⇒ check the encirclement around the point $(-1/K, 0)$

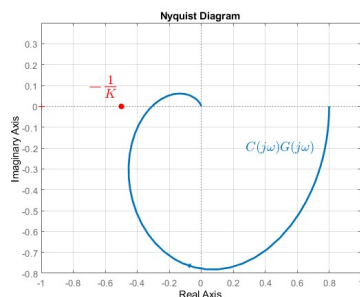
Example with $C(s)G(s) = \frac{1}{(s+1.5)(s-1)}$ and $K = 2$



Actually system
stable
 $\forall K > 1.5$

Another Example with $C(s)G(s) = \frac{1}{s^3 + 2s^2 + 2.25s + 1.25}$ and $K = 2$

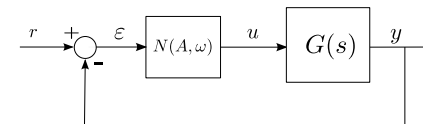
(poles = -1 and $-0.5 \pm j$)



Actually system stable for $-1.25 < K < 3.25$

Stability of limit cycles

Back to closed-loop system with a describing function :

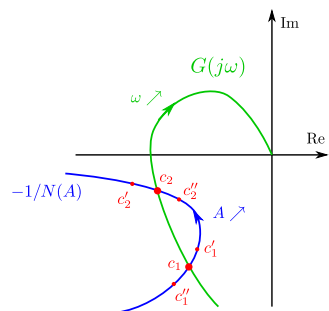


Characteristic equation : $1 + N(A, \omega)G(s) = 0$ or $G(s) = -\frac{1}{N(A, \omega)}$

- ▶ Check encirclement around the point $(\text{Re}[-\frac{1}{N}], \text{Im}[-\frac{1}{N}])$
- ▶ By assumption $G(s)$ is stable \rightarrow no unstable pole
- ▶ Check if critical points is on left or right of $G(j\omega)$ locus when $\omega \nearrow$

Graphical analysis

Consider a system where two limit cycles are predicted



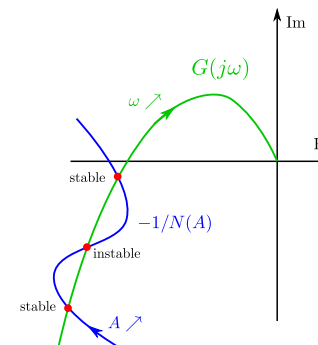
- ▶ At c_1 , there is a limit cycle with an amplitude A_1 and a frequency ω_1
 - if a slight disturbance increases A ; we move to c_1' ; the system is unstable; the amplitude continues to increase; we move along curve $-1/N(A)$ toward c_2
 - if a slight disturbance decreases A ; we move to c_1'' ; the system is stable; the amplitude continues to decrease; we move along curve $-1/N(A)$ toward 0
 - the limit cycle is unstable
- ▶ At c_2 , there is a limit cycle with an amplitude A_2 and a frequency ω_2
 - if a slight disturbance increases A ; we move to c_2' ; the system is stable; the amplitude decreases; we move back toward c_2
 - if a slight disturbance decreases A ; we move to c_2'' ; the system is unstable; the amplitude increases; we move back toward c_2
 - the limit cycle is stable

Stability condition (graphical)

Loeb criterion

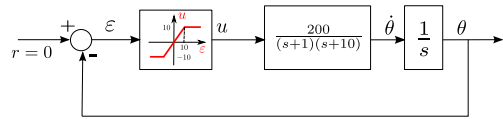
A limit cycle of amplitude A_0 and frequency ω_0 is stable if the intersection point is such that along the nyquist plot of $G(j\omega)$ as ω increases, the direction of increasing A along the critical curve $-1/N(A)$ is toward the left.

Example :



Example

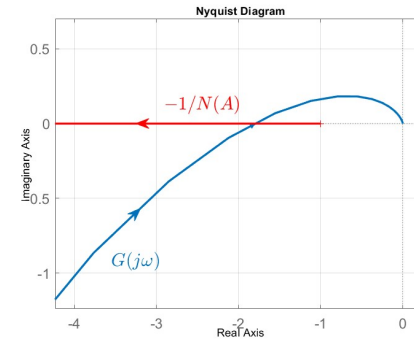
Simple control of a DC motor with a saturation



- ▶ What is the describing function of the nonlinearity?
- ▶ Show that a limit cycle exists.
- ▶ What would be the approximated amplitude and frequency of the self-oscillations?
- ▶ Is the limit cycle stable?

Example

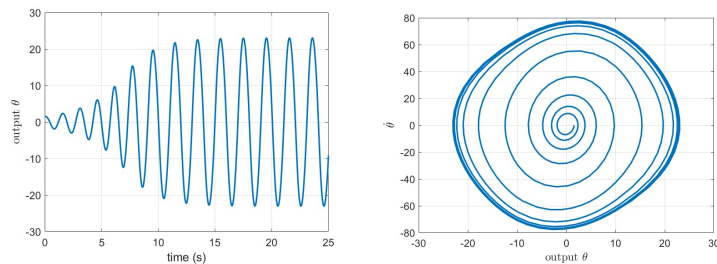
- ▶ Describing function : $N(A) = \frac{2}{\pi} \arcsin\left(\frac{10}{A}\right) + \frac{20}{A\pi} \sqrt{1 - \frac{100}{A^2}}$ (if $A > 10$)
- ▶ Nyquist plot of $G(j\omega)$ and $-1/N(A)$; note that $N(10) = 1$ and $N(+\infty) = 0$



- ▶ From the plot, $\omega = 3.22 \text{ rad/s}$ ($T = 1.95 \text{ s}$) and $A = 22.1$
- ▶ The limit cycle is stable

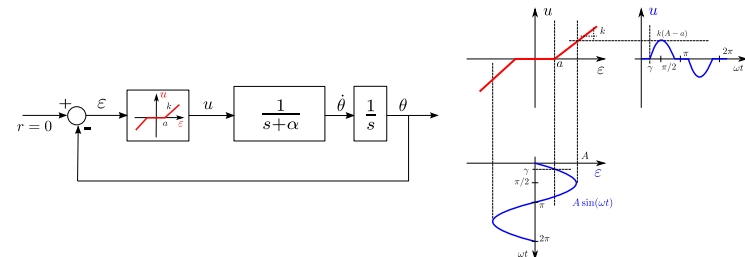
Example

Simulation of the closed loop : output (left) and phase plane (right)



Exercise

Simple control of a DC motor with a dead zone



- ▶ What is the describing function of the nonlinearity?
- ▶ Draw a sketch of the nyquist plot of $G(j\omega)$ and the critical locus $-1/N(A)$.
- ▶ Does a limit cycle exists?

Solution